

Terahertz calculations for the Australian synchrotron

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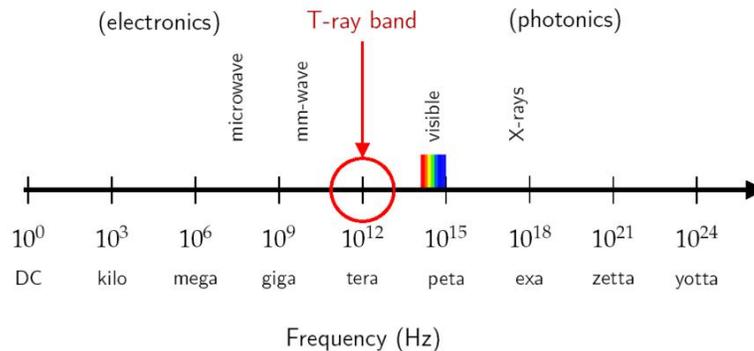
ABSTRACT

The development of terahertz radiation (T-rays) is spurring new applications in spectroscopy and imaging. To maximize the use of T-rays in more applications, a high average terahertz power is needed. Rather than using fast diodes or laser sources, this paper will show that a synchrotron can generate high average power T-rays. This is achieved by creating an electron bunch in the synchrotron ring with high intensity in the terahertz frequency region via Thomson scattering.

Keywords: Synchrotron, Terahertz, T-rays

1. INTRODUCTION

T-rays are a form of electromagnetic radiation, lying between millimeter waves and infrared radiation on the electromagnetic spectrum, as shown in the diagram below. T-ray radiation interacts in a unique way with matter, relating to rotational and vibration modes in gas, inter-domain movement in large biomolecules, and resonances in crystals. T-rays have been used to identify water levels in living tissue, measure the real-time 'dressing' of electrons in semiconductors^{1,2,3}, detect malignant skin cancer^{4,5}, and detect Anthrax-like spores in envelopes⁶. A newly accessible region of the electromagnetic spectrum will provide new technologies, especially if the low-power limitations of T-ray systems can be overcome.



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In the past decade, broadband T-ray sources have been developed using fast diodes, femtosecond lasers and non-linear transmission lines, leading to increased interest in T-ray applications. However, the low average power generated by these methods⁷ (about 10^{-6} W) is a major hurdle to be overcome. For example, the emerging technology of terahertz time domain spectroscopy (THz-TDS) provides the ability to identify a substance via its absorption spectra. With higher average power T-ray radiation, better detection of concealed weapons, hidden explosives and land mines can be performed.

In November 2002, the Jefferson Laboratory in the USA produced high average power broadband emission from subpicosecond electron bunches in an accelerator by using the facility's Energy Recovery Linac (ERL). The average power obtained was nearly 20 Watts, many times larger than other T-ray sources⁷. Jefferson Laboratory's ERL is a superconducting radio-frequency electron accelerator. The ERL is designed to recover spent by the electron bunches. In the ERL, short bunches of electrons, each with a pulse width of about 500 femtoseconds, were accelerated to energies of about 40 MeV and then accelerated through a bend using a powerful magnetic field. The acceleration of the electron bunches through this bend caused T-ray radiation to be generated with a peak power⁷ of 106 W and a peak frequency⁷ of 0.6 THz.

This paper models T-ray generation in a synchrotron ring. In the linear accelerator (linac) system that Jefferson Laboratory used, each pass of the electron beam through the magnetic bend carried new electrons. However, in a synchrotron ring, the same electrons pass through the magnetic bend repeatedly without the replacing of new electrons.

In this paper, terahertz frequencies are theoretically obtained by narrowing the electron bunch via Thompson scattering. As scattering can be achieved with an ultrafast laser, gated detection via the standard pump-probe method^{3,8,9} can be used. The pump-probe approach has the added benefit of reducing jitter.

2. SYNCHROTRON RADIATION

A synchrotron ring emits a broadband spectrum of light by accelerating bunches of electrons at relativistic speeds in the ring. The intensity of the synchrotron radiation can be quantified by its brightness and power. Power can be represent by units of either Watts/cm⁻¹ or photons/sec/0.1%bandwidth. However, brightness or brilliance is defined as power divided by emittance. Therefore, the brightness would have units of either Watts/cm⁻¹/mm²/sr or photons/sec/0.1%bandwidth/mm² sr.

2.1. Power

The power emitted from a synchrotron ring can be calculated by using the following formula:

$$P(\lambda) = 4.38 \cdot 10^{14} I \cdot \theta \cdot BW \cdot \left(\frac{\rho}{\lambda}\right)^{1/3} \text{ photons/sec} \quad (1)$$

where I is the current in amperes, θ is the horizontal collection angle in radians, BW is the bandwidth in %, λ is the wavelength in microns and ρ is the radius of the ring and it has the same units as λ . Equation (1) shows the power of the synchrotron radiation in photons/sec, but it can be converted to Watts by dividing P by $(5.04 \times 10^{18} \times \lambda)$. For a current (I) of 200 milliamperes and a 0.1% bandwidth (BW), Equation (1) can be simplified as:

$$P(\lambda) = \frac{0.1\lambda \cdot 2(f_{pa}(\lambda) + f_{pe}(\lambda))}{5.04 \cdot 10^{18} \lambda} \text{ Watts/cm}^{-1} \quad (2)$$

where the 0.1λ term results from the 0.1% bandwidth and therefore,

$$P(\lambda) = 3.9682 \cdot 10^{-20} (f_{pa}(\lambda) + f_{pe}(\lambda)) \text{ Watts/cm}^{-1}. \quad (3)$$

The parameters $f_{pa}(\lambda)$ and $f_{pe}(\lambda)$ are the parallel and perpendicular polarized flux per milliradian vertically, which depend on the value of λ in microns.

2.2. Flux

The formulae for calculating the values of $f_{pa}(\lambda)$ and $f_{pe}(\lambda)$ are¹⁰:

$$f_{pe}(\lambda) = 1.3256 \cdot 10^{16} E^2 \theta_h(\lambda) k_1(\lambda) \quad (4)$$

$$f_{pa}(\lambda) = 1.3256 \cdot 10^{16} E^2 \theta_h(\lambda) k_2(\lambda). \quad (5)$$

Here, E is the synchrotron beam energy in units of electron volts, while $\theta_h(\lambda)$, $k_1(\lambda)$ and $k_2(\lambda)$ are the horizontal emission angles in radians and the intermediate variables, which depend on the values of λ , are given as:

$$k_1(\lambda) = 2.5 \cdot 10^{-10} \lambda^2 \gamma^2 \theta_h(\lambda)^2 \cdot (1 + 10^{-6} \gamma^2 \theta_h(\lambda)^2) \cdot (\exp(\frac{-x_1(\lambda)}{2}) + G_1(\lambda))^2 \quad (6)$$

$$k_2(\lambda) = 5 \cdot 10^{-4} \lambda^2 \cdot [(1 + 10^{-3} \gamma \theta_h(\lambda))^2 \cdot (\exp(\frac{-x_1(\lambda)}{2}) + G_2(\lambda))]^2 \quad (7)$$

where γ is the electron relativistic to rest mass ratio, which is equal to $1957 \times E$ and E is the synchrotron beam energy in GeV. Also $x(\lambda)$ is the ratio of critical wavelength over wavelength (λ_c/λ). By definition, the critical wavelength is the wavelength which bisects the power spectrum and it is defined as $(5.59 \times \rho)$ (meters) / E^3 (GeV). Consequently, we can define $x(\lambda)$ as:

$$x(\lambda) = \frac{5.59 \rho}{10^4 E^3 \lambda}. \quad (8)$$

2.3. Bessel functions

In Equations (6) and (7), $G_1(\lambda)$ and $G_2(\lambda)$ are Bessel functions and $x_1(\lambda)$ is also used in obtaining the required Bessel functions and is defined as¹⁰:

$$x_1(\lambda) = \frac{1}{2} x(\lambda) (1 + (10^{-6} \lambda^2 \theta^2)^{3/2}). \quad (9)$$

Note that $G_1(\lambda)$ is the Bessel function of order 1/3, for light polarized perpendicularly to the orbit plane. However, $G_2(\lambda)$ is the Bessel function of order 2/3, for light, polarized parallel to the orbit plane. By definition, the Bessel function will be an expression of the form:

$$G(x) = x \int_x^\infty K_{5/3}(\eta) d(\eta). \quad (10)$$

By knowing that x is the ratio of the critical wavelength over wavelength, the integral over the spherical fractional order Bessel function in Equation (10) can now be expressed using the Kostroun algorithm^{11,12}, with fractional order, k :

$$G_k(x) = \frac{1}{2} \left(\frac{e^{-x}}{2} + \sum_{r=1}^{\infty} e^{-x \cdot \cosh(0.5r)} \cdot \cosh(0.5kr) \right). \quad (11)$$

2.4. Brightness

As discussed earlier, brightness is defined as the ratio of power by emittance ($P(\lambda)/\epsilon$). By definition, emittance is a measurement of the synchrotron radiation intensity lost due to divergence of the electron beam and its non-zero cross-sectional size. However, in practical terms, the brightness is the flux divided by the opening angle and the area of the source. As a consequence, low emittance will result in high brightness of synchrotron radiation. The emittance, ϵ can be expressed by using the following equation:

$$\epsilon = 10^{-12} \theta_h \theta_v S_h S_v \text{ mm}^2 \text{ sr.} \quad (12)$$

The θ_h and θ_v terms are the horizontal and vertical emission angles in milliradians while S_h and S_v are the horizontal and vertical source sizes in microns as shown in Figure 1. The horizontal and vertical source sizes can be obtained by using the equations:

$$S_h = \sqrt{S_{hi}^2 + S_{hp}^2 + S_{hdiff}^2}$$

$$S_v = \sqrt{S_{vi}^2 + S_{vp}^2 + S_{vdiff}^2} \quad (13)$$

where S_{hi} and S_{vi} are the intrinsic horizontal and vertical source sizes, S_{hp} and S_{vp} are the source sizes due to projection and S_{hdiff} and S_{vdiff} are the source sizes due to diffraction.

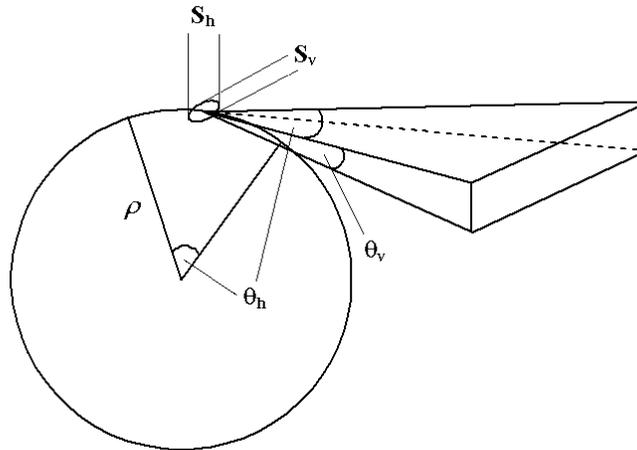


Figure 1. Synchrotron parameters. The circle in the diagram is the bending circle of a synchrotron ring with radius, ρ . The S_h and S_v are horizontal and vertical source sizes in microns and the θ_h and θ_v are the horizontal and vertical emission angles in milliradians.

To calculate the horizontal and vertical emission angle (θ_h and θ_v), the natural opening angle θ_{nat} and the beamline vertical and horizontal angle (θ_{hb} and θ_{vb}) are needed. Once these angles are determined, the horizontal emission angle will be the smallest angle of either θ_{nat} or θ_{hb} and the vertical emission angle will be the smallest angle between θ_{nat} and θ_{vb} . In other words, the θ_h and θ_v must be less than or equal to θ_{nat} . The natural opening angle, θ_{nat} can be determined by using Equation (14) where λ is in microns and ρ is in meters¹⁰,

$$\theta_{nat} = 1.66188 \cdot \left(\frac{10^3 \lambda}{\rho} \right)^{1/3}. \quad (14)$$

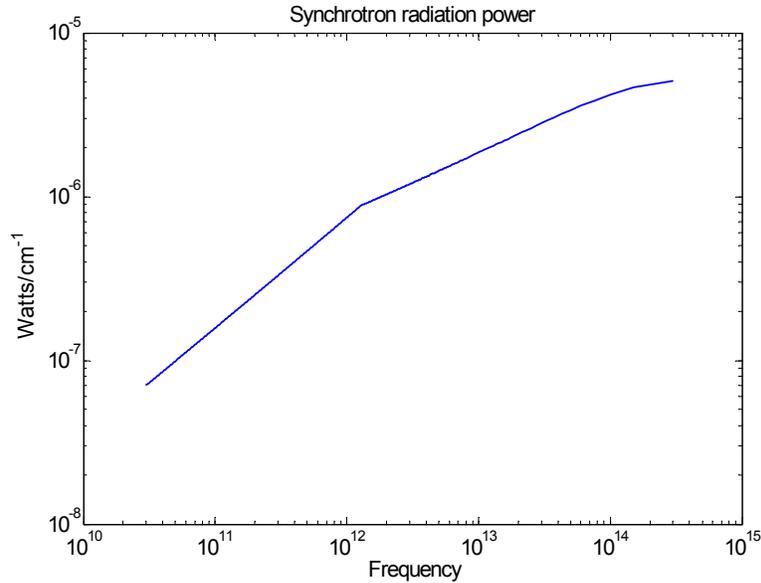


Figure 2. Synchrotron radiation power. For a current of 0.2 ampere, bandwidth of 0.1%, sources size of 500 microns (horizontal) and 200 microns (vertical) and a 60 mradian (horizontal) by 60 mradian (vertical) collection angles, Equation (3) is used to generate the graph with Watts/cm⁻¹ versus frequency where cm⁻¹ represents the wavenumber. Synchrotron radiation power is higher for increased spectral frequency. At a frequency of one terahertz (10¹² Hz), the intensity is around 10⁻⁶ Watts/cm⁻¹. This demonstrates that only weak T-ray power is generated in a regular synchrotron.

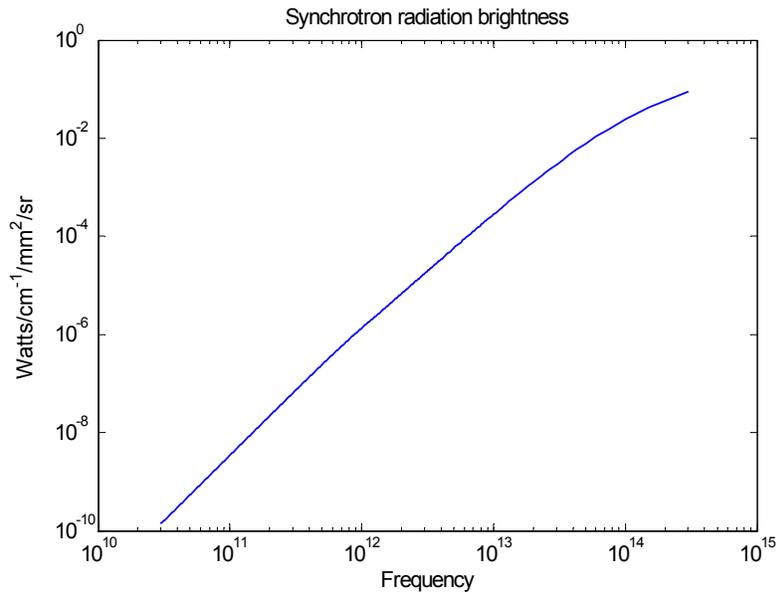


Figure 3. Synchrotron radiation brightness. By definition, the brightness equals to the power divided by emittance (P/ϵ). Therefore, with the same parameters used in Figure 2, the synchrotron brightness curve has the same characteristic as the synchrotron power curve. Similar to Figure 2, the synchrotron brightness intensity will increase when the frequency increases and this also shows a weak intensity occurs in the terahertz frequency range.

Figure 2 and Figure 3 show the synchrotron radiation power and brightness plot against spectral frequency. Moreover, it can be clearly seen that the synchrotron radiation is weak in the terahertz (10^{12} Hz) frequency region. To obtain high intensity synchrotron radiation, Thomson scattering is needed.

3. THOMSON SCATTERING FOR T-RAY GENERATION

The spectra of Figures 2 and 3 were derived from the theory of the passage of a single electron. In practice many electrons circulate together in bunches, adding another time-scale to the calculation. The squared Fourier transform of a Gaussian pulse will yield the intensity of the synchrotron radiation as shown in Figures 2 and 3. If the electron Gaussian pulse is sliced to have a shorter duration, the Fourier transform of the time domain characteristic will give a stronger intensity in terahertz region. If the electron bunch's dimensions are small compared to the wavelength of observation, the Fourier transform of the electron bunch will reach the terahertz region after slicing. The technique that is used to slice the electron bunch is known as *locked laser Thomson scattering*. A 500 MHz mode-locked laser will shoot into a certain part of the synchrotron ring and slice the electron bunches to become narrower. When the size-reduced electron bunches are viewed in the frequency domain, an enhancement occurs in the frequency domain plot and the frequency reaches the terahertz region. The coherent enhancement factor, N of Thomson scattering is given as:

$$N = 1 - f(\omega) + f_s n_e f(\omega) \quad (15)$$

where n_e is the number of electrons per bunch, f_s is the slicing factor. If the electron bunch is sliced by 200 femtoseconds every 20 picoseconds, the slicing factor f_s will be 0.01. Here, $f(\omega)$ represents the Fourier transform calculation, which is given by¹⁰:

$$f(\omega) = \exp\left(\frac{-4\pi^2\sigma^2}{\lambda^2}\right) \quad (16)$$

where λ represents the wavelength in microns and σ is the bunch sigma value in microns. The bunch σ is given by:

$$\sigma = \frac{0.3 L_{fwhm}}{2.2} \quad (17)$$

where the L_{fwhm} is the bunch length full width half maximum value in femtoseconds. When the coherent enhancement factor has been determined, the new power spectrum of the synchrotron radiation from Equation (3) is now:

$$P'(\lambda) = P(\lambda) \cdot N. \quad (18)$$

4. MATLAB SIMULATION RESULT

From Figure 4 and Figure 5, it can be seen that the synchrotron radiation power and brightness graphs occur to have an enhancement area. This enhancement area is not present before Thomson scattering as we can see by comparing Figures 4 and 5 with Figures 2 and 3. These enhancement areas occur due to the effect of narrowing the electron bunches via Thomson scattering. When the slices electron bunches become thinner and smaller, the power spectra of the electron bunch will result a high intensity in terahertz frequency region.

From Figure 4, a power of about $13.5 \text{ Watts/cm}^{-1}$ is generated at 10^{12} Hz. However, the peak in Figure 5 does not occur in the terahertz frequency region. A brightness of about $25 \text{ Watts/cm}^{-1}/\text{mm}^2/\text{sr}$ is generated but this is much higher than the existing one in Figure 3, which is around $1.5 \cdot 10^{-6} \text{ Watts/cm}^{-1}/\text{mm}^2/\text{sr}$. As a consequence, this demonstrates that the Thomson scattering slicing method is an efficient way of producing T-ray power using synchrotron radiation.

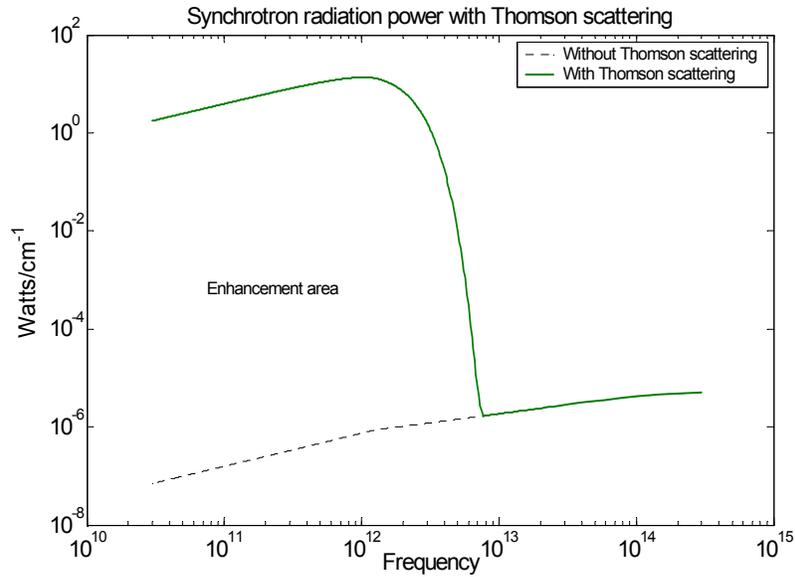


Figure 4. Synchrotron radiation power with Thomson scattering. Using Equation (18), with $f_s= 0.01$, an enhancement area occurs in the graph. The graph with a solid line represents the synchrotron radiation power after Thomson scattering has been applied and the dotted line represents the synchrotron power before Thomson scattering. The graph also shows a peak occurs in the terahertz frequency region.

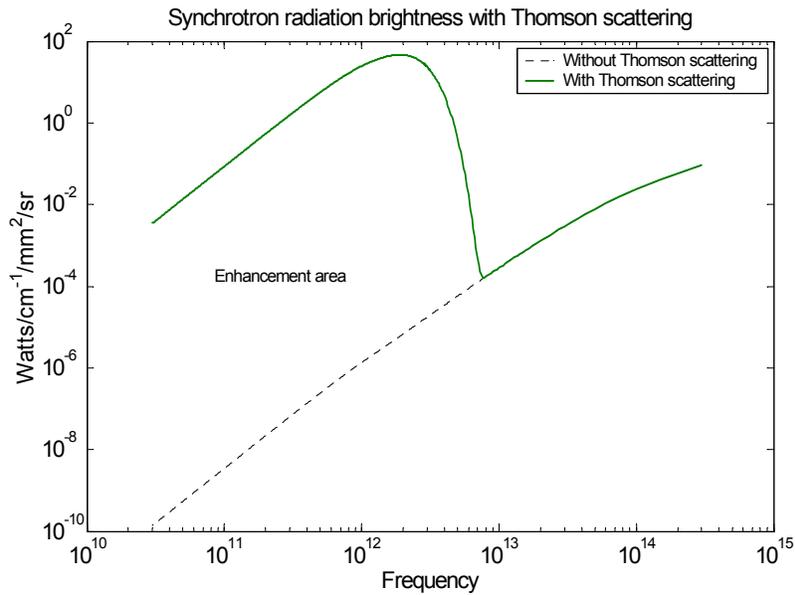


Figure 5. Synchrotron radiation brightness with and without Thomson Scattering. Using the same coherent enhancement factor, N multiplied by the brightness in Figure 3, with high brightness intensity is generated as demonstrated in the above graph. The graph shows an enhancement similar to that in Figure 4 and is just a different way of representing the synchrotron radiation intensity.

5. CONCLUSION

We have shown the calculations required for estimating T-ray power output from a synchrotron using Thomson scattering. Moreover, high synchrotron T-ray power will maximize the application of T-rays in several fields. Note that the calculated power $P(\lambda)$ depends on many parameters such as the bending radius, beamline vertical and horizontal angle, bunch length, source sizes, etc. In practice, many parts of the synchrotron ring need to be considered such as the bending magnet, booster ring, linear accelerator electron gun, and storage ring situation.

ACKNOWLEDGEMENTS

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REFERENCES

1. D. Grischkowsky, S. Keiding, M. V. Exter, C. Fattinger, "Far-infrared time-domain spectroscopy with terahertz beams of dielectrics and semiconductors," *Journal of the Optical Society of America B: Optical Physics* **7**(10), pp. 2006-2015, 1990.
2. D. M. Mittleman, J. Cunningham, M. C. Nuss, M. Geva, "Noncontact semiconductor wafer characterization with the terahertz Hall effect," *Applied Physics Letters* **71**(1), pp. 16-18, 1997.
3. S. P. Mickan, D. Abbott, J. Munch, X. C. Zhang, and T. van Doorn, "Analysis of system trade-offs for terahertz imaging," *Microelectronics Journal* (Elsevier) **31**(7), pp. 503-514, 2000.
4. R. M. Woodward, B. E. Cole, V. P. Wallace, D. D. Arnone, R. J. Pye, E. H. Linfield, M. Pepper, "Terahertz pulse imaging of in reflection geometry of human skin cancer and skin tissue," *Physics in Medicine and Biology* **47**, pp. 3853-3863, 2002.
5. B. Ferguson, S. Wang, D. Gray, D. Abbott, X.-C. Zhang, "Identification of biological tissue using chirped probe THz imaging," *Microelectronics Journal* (Elsevier) **33**, pp. 1043-1051, 2002.
6. S. Wang, B. Ferguson, C. Mannella, D. Gray, D. Abbott, X.-C. Zhang, "Powder detection using THz imaging," *OSA Trends in Optics and Photonics (TOPS)* **73**, pp. 132, 2002.
7. G. L. Carr, Michael C. Martin, Wayne R. McKinney, K. Jordan, George R. Neil, and G. P. Williams, "High-power terahertz radiation from relativistic electrons," *Nature* **420**, pp.153-156, 2002.
8. S. P. Mickan, K. S. Lee, T. M. Lu, J. Munch, D. Abbott, and X.-C. Zhang, "Double modulated differential THz-TDS for thin film dielectric characterization," *Microelectronics Journal* (Elsevier) **33**(12), pp. 1033-1042, 2002.
9. S. P. Mickan, A. Menikh, H. Liu, C. A. Mannella, D. Abbott, J. Munch, X.-C. Zhang, "Label-free bioaffinity detection using terahertz technology," *Physics in Medicine and Biology* **47**(21), pp. 3789-3795, 2002.
10. G. P. Williams, "FAR-IR/THz radiation from the Jefferson Laboratory, energy recovered linac, free electron laser," *Review of Scientific Instruments* **73**(3), pp. 1461-1463, 2002.
11. V. Kostroun, "Simple numerical evaluation of modified Bessel functions $K_\nu(x)$ of fractional order and the integral $x^\nu K_\nu(\eta)d(\eta)$," *Nuclear Instruments & Methods* (Elsevier) **172**, pp. 371-374, 1980.
12. H. N. Rutt, "On the evaluation of modified Bessel functions of the second kind and fractional order for synchrotron radiation calculations," *Nuclear Instruments & Methods in Physics Research Section A* (Elsevier), 2003 (In Press).

APPENDIX

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% MATLAB CODE
% Synchrotron parameters calculation for terahertz (T-ray) generation
clear all;
E = 3.0; % Synchrotron beam energy in GeV
rho = 5; % Synchrotron bending radius in meters
thetavb = 60; % Beamline vertical angle in milliradians
thetahb = 60; % Beamline horizontal angle in milliradians
shi = 500; % Intrinsic horizontal source size in microns
svi = 200; % Instinsic vertical source size in microns

cpb = 400; % Charge per bunch in pico-coulombs
etendue = 1*10.^(-4); % Etendue (throughput) in mm^2 steradians use 1e-4 for 10 microns pinhole
fwhm = 200; % Bunch length FWHM value in femtoseconds here
sigmab = (3*10.^14 * fwhm * (1*10.^(-15))) / 2.2; % Bunch sigma value in microns
ne = (cpb * (1*10.^(-12))) / (1.6*10.^(-19)); % Number of electrons/bunch

lambda = 1:1:10000; % Wavelength
omega = (lambda * 0.0001).^(-1);
domega = 1;
factor = 1000 * domega ./ omega;

% Calculate the horizontal and vertical emission angle
for i=1:1:10000
    thetavnat(i) = 1.66188 * ((1000 * lambda(i)) / rho).^0.33333;
    if thetavnat(i) < thetavb
        thetav(i) = thetavnat(i);
    else
        thetav(i) = thetavb;
    end

    if thetavnat(i) < thetahb
        thetah(i) = thetavnat(i);
    else
        thetah(i) = thetahb;
    end
end

% Calculate the source sizes due to projection and diffraction
for m=1:1:10000
    shp(m) = rho * thetah(m) * thetah(m) / 8;
    shdiff(m) = lambda(m) / (thetah(m) / 1000);
    svp(m) = rho * thetav(m) * thetah(m) / 8;
    svdiff(m) = lambda(m) / (thetav(m) / 1000);
    % Calculate the horizontal and vertical source sizes
    sh(m) = ((shi.^2) + (shp(m).^2) + (shdiff(m).^2)).^0.5;
    sv(m) = ((svi.^2) + (svp(m).^2) + (svdiff(m).^2)).^0.5;
    emittance(m) = thetah(m) * thetav(m) * sh(m) * sv(m) * (1*10.^(-12));
end

sum = 0;
h = 0.5;
n1 = 1/3;

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n2 = 2/3;
for n=1:1:10000
x(n) = (5.59 * rho / (E.^3)) / (lambda(n) * 10000);
end
gamma = 1957 * E;

% Calculate the Bessel Function G1
for p=1:1:10000
sum(p) = 0;
for theta = 0.5:0.25:(thetav/2)
sum1(p) = 0;
x1(p) = (x(p)/2) * ((1 + (gamma * 0.001 * theta).^2).^(3/2));
for q=1:50
cosh(q) = (exp(q * h * n1) + exp((-q) * h * n1)) / 2;
csh(q) = (exp(q * h) + exp((-q) * h)) / 2;
gi(q) = (exp((-x1(p)) * csh(q)) * (cosh(q)));

if gi(q) > 0.00001
sum1(p) = sum1(p) + gi(q);
end
end

sum2(p) = 0;
x1(p) = (x(p)/2) * ((1 + (gamma * 0.001 * theta).^2).^(3/2));
for q=1:50
cosh(q) = (exp(q * h * n2) + exp((-q) * h * n2)) / 2;
csh(q) = (exp(q * h) + exp((-q) * h)) / 2;
gi(q) = (exp((-x1(p)) * csh(q)) * (cosh(q)));

if gi(q) > 0.00001
sum2(p) = sum2(p) + gi(q);
end
end

% Calculate the flux polarized parallel (fpa) or perpendicular (fpe) per milliradians vertically, in other words from
% (theta-0.5)<theta<(theta+0.5) in milliradians. To do this, we multiply k1 and k2 by 0.001 radians
k1(p) = 0.001 * x(p).^2 * ((gamma * 0.001 * theta).^2) * (1 + (gamma * 0.001 * theta).^2) * ((h * ((exp(-(x1(p)))) /
2) + sum1(p))).^2);
k2(p) = 0.001 * x(p).^2 * ((1 + (gamma * 0.001 * theta).^2) * h * ((exp(-(x1(p)))) / 2) + sum2(p)).^2;
fpe(p) = (1.3256 * 10.^16) * E.^2 * thetah(p) * k1(p);
fpa(p) = (1.3256 * 10.^16) * E.^2 * thetah(p) * k2(p);
sum(p) = sum(p) + 2 * (fpa(p) + fpe(p));
end

end

%sum gives the result in photons/sec into 0.1% bandwidth using the smaller
%of the natural and actual vertical opening angles, and the smaller of the
%natural and actual horizontal angles, but the latter can be changed
%if we need Watts per cm-1, then multiply by the bw factor and convert
for l=1:1:10000
watts(l) = factor(l) * sum(l) / ((5.04 * 10.^24) * lambda(l) * 0.000001);
%we have total flux, now calculate the brightness
bright(l) = sum(l) / emittance(l);

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wattsbrt(l) = watts(l) / emittance(l);
if etendue > emittance(l)
    sampsig(l) = watts(l);
else
    sampsig(l) = watts(l) * etendue / emittance(l);
end
end

%calculate of the sample signal, change sampsig to sum of photons/sec
%calculate f(w)
for a=1:1:10000
    if lambda(a) < 40
        fofw(a) = 0;
    else
        fofw(a) = exp(-((4*(3.14.^2)*(sigmab.^2)) / lambda(a).^2));
    end
    % Calculate the coherent enhancement factor
    cef(a) = (1 - fofw(a)) + (0.01 * ne * fofw(a));
    cohwt(a) = watts(a) * cef(a);
    cohwtbrt(a) = wattsbrt(a) * cef(a);
end

figure(1)
loglog((3*10.^8)./(lambda.*10.^(-6)),watts);
title('Synchrotron radiation power');
xlabel('Frequency');
ylabel('Watts/cm^{-1}');

figure(2)
loglog((3*10.^8)./(lambda.*10.^(-6)),wattsbrt./5);
title('Synchrotron radiation brightness');
xlabel('Frequency');
ylabel('Watts/cm^{-1}/mm^2/sr');

figure(3)
loglog((3*10.^8)./(lambda.*10.^(-6)),watts,(3*10.^8)./(lambda.*10.^(-6)),cohwt);
title('Synchrotron radiation power with Thomson scattering');
xlabel('Frequency');
ylabel('Watts/cm^{-1}');

figure(4)
loglog((3*10.^8)./(lambda.*10.^(-6)),wattsbrt./5,(3*10.^8)./(lambda.*10.^(-6)),cohwtbrt./5);
title('Synchrotron radiation brightness with Thomson scattering');
xlabel('Frequency');
ylabel('Watts/cm^{-1}/mm^2/sr');

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